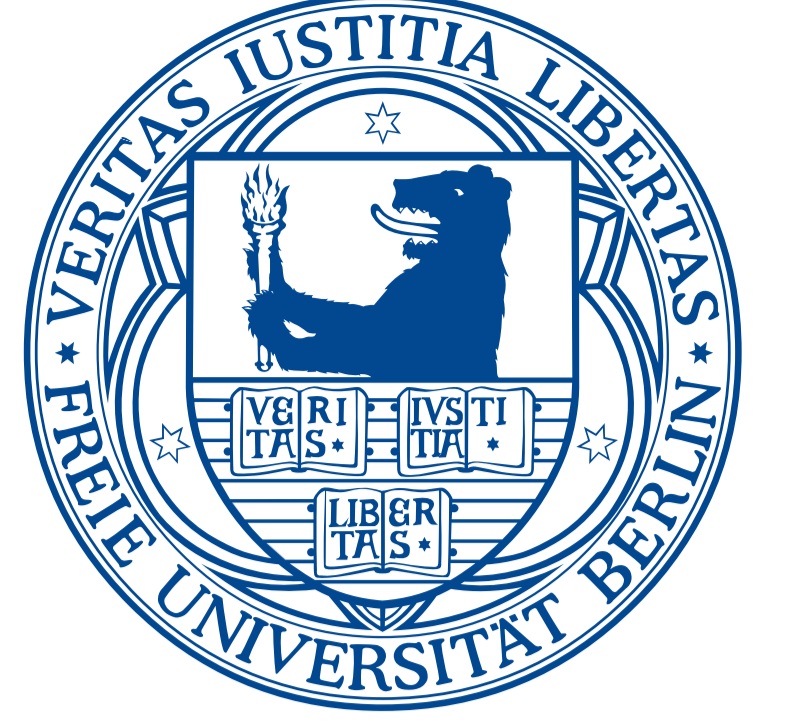




# A sharp interface finite volume method for variable density zero Mach number two-phase flow with soluble surfactants

Stephan Gerber<sup>†</sup>, Matthias Waidmann<sup>‡</sup>  
Michael Oevermann<sup>†</sup>, Rupert Klein<sup>‡</sup>

<sup>†</sup>Konrad-Zuse-Zentrum für Informationstechnik Berlin  
<sup>‡</sup>Freie Universität Berlin, Institut für Mathematik, AG Geophysical Fluid Dynamics



## Zero-Mach number variable density finite volume scheme

Leading order system in dimensional (space-)integral form

$$\begin{aligned} \text{mass: } & \frac{d}{dt} \int_{\Omega} \rho dV = - \int_{\partial\Omega} (\rho \theta \mathbf{u} \cdot \mathbf{n}) \left( \frac{1}{\theta} \right) dS \\ \text{spec: } & \frac{d}{dt} \int_{\Omega} \rho Y_s dV = - \int_{\partial\Omega} (\rho \theta \mathbf{u} \cdot \mathbf{n}) \left( \frac{Y_s}{\theta} \right) dS - \int_{\partial\Omega} \mathbf{j}_s \cdot \mathbf{n} dS + \int_{\Omega} \dot{\sigma}_s dV \\ \text{mom: } & \frac{d}{dt} \int_{\Omega} \rho \mathbf{u} dV = - \int_{\partial\Omega} (\rho \theta \mathbf{u} \cdot \mathbf{n}) \left( \frac{\mathbf{u}}{\theta} \right) dS - \int_{\partial\Omega} p^{(2)} \mathbf{n} dS + \int_{\Omega} \rho \mathbf{g} dV + \int_{\partial\Omega} \mathbf{T} \cdot \mathbf{n} dS + \int_{\partial\Omega \cap \Gamma} \sigma \mathbf{t} d\Gamma \\ \text{en: } & \frac{d}{dt} \int_{\Omega} \rho \theta dV = - \int_{\partial\Omega} (\rho \theta \mathbf{u} \cdot \mathbf{n}) dS + \int_{\Omega} \rho \theta D dV \quad \text{with} \\ & \int_{\Omega} \rho \theta D dV = - \frac{dp^{(0)}}{dt} \int_{\Omega} \frac{\theta}{c^2} dV + \int_{\Omega} \left( \frac{\theta}{c^2} \left[ \nabla \cdot \mathbf{q} + \sum_s \mathbf{j}_s \cdot \nabla h_s + \sum_s \Delta h_s \dot{\sigma}_s \right] \right) dV \end{aligned}$$

*black:* implemented as predictor (2<sup>nd</sup> order Runge-Kutta) - corrector (2 projection steps) method (yet single-phase) with fluxes based on fluxes of  $\rho \theta$  where  $\theta$  is the entropy  
*red:* not implemented yet

1<sup>st</sup> projection step: Poisson equation to enforce divergence constraint via flux corrections

$$\int_{\partial\Omega} \left[ \frac{\Delta t}{2} \theta^{(n+\frac{1}{2})} \right] (\nabla \pi^{(2)})^{(n+\frac{1}{2})} \cdot \mathbf{n} dS = - \left( \frac{d}{dt} \int_{\Omega} \rho \theta dV \right)^{(*)} \quad \text{where } (\nabla \pi^{(2)})^{(n+\frac{1}{2})} = (\nabla \pi^{(2)})^{(n+\frac{1}{2})}$$

2<sup>nd</sup> projection step: Poisson equation to correct pressure via momentum flux correction

$$\int_{\partial\Omega} \left[ \Delta t \theta^{(n+1)} \right] (\nabla \partial p^{(2)})^{(n+\frac{1}{2})} \cdot \mathbf{n} dS = \int_{\partial\Omega} (\rho \theta \mathbf{u})^{(n+1,**)}. \mathbf{n} dS - \int_{\Omega} (\rho \theta D)^{(n+1,**)}. dV$$

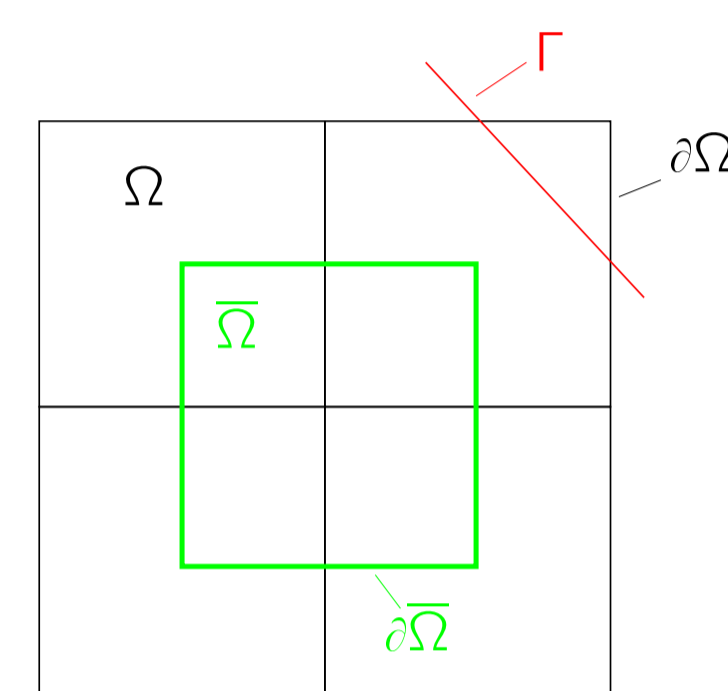


Fig 1: interface (red), primal (black) and dual (green) cells

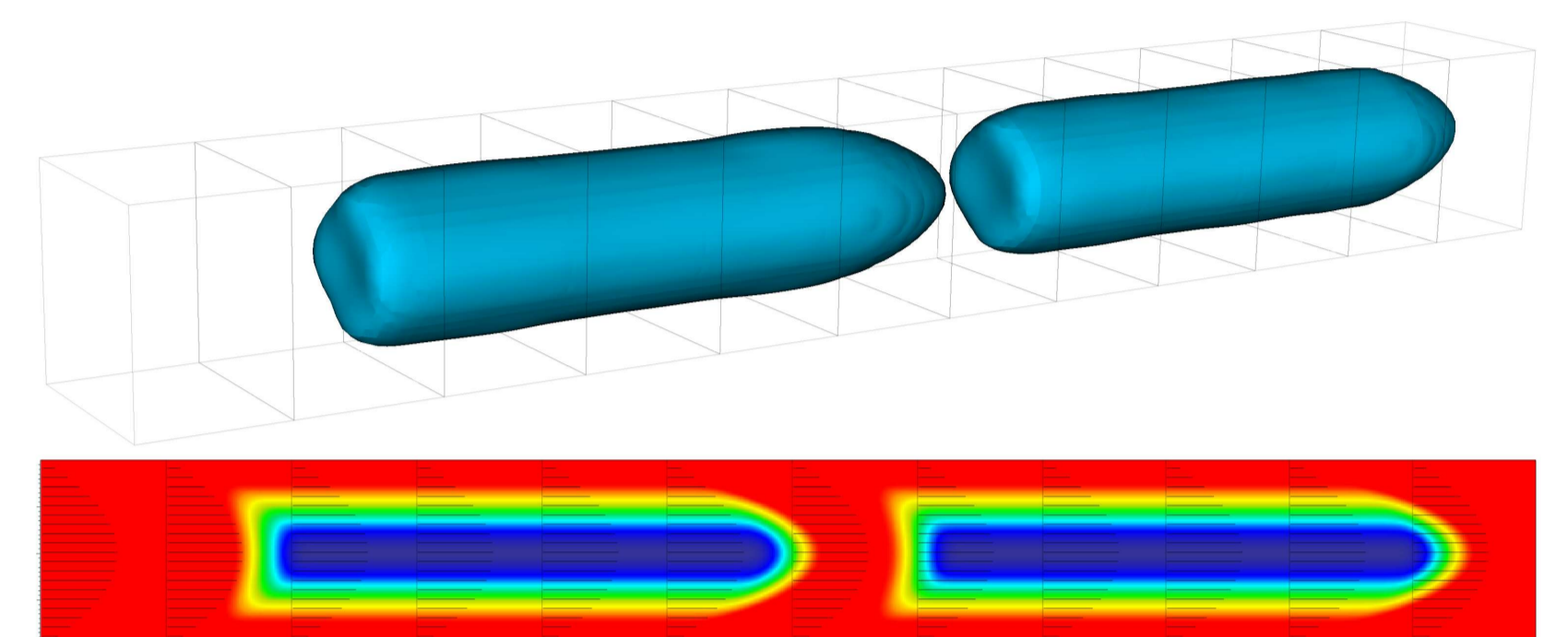


Fig 2: bubbles (density) in horizontal channel at zero gravity (yet no interface / phase separation included)

## Discretisation of a surface surfactant equation with an embedding method\*

Prototype of an advection-diffusion equation on an interface  $\Gamma \subset \mathbb{R}^n$ :

$$\frac{\partial \phi_{\Gamma}}{\partial t} = F(\phi_{\Gamma}, \nabla_{\Gamma} \phi_{\Gamma}, \Delta_{\Gamma} \phi_{\Gamma}), \quad (1)$$

with solution  $\phi_{\Gamma} \in \mathbb{R}$ , surface gradient  $\nabla_{\Gamma}$  and Laplace-Beltrami operator  $\Delta_{\Gamma} = \nabla_{\Gamma} \cdot \nabla_{\Gamma}$ .

Instead of solving (1) on the surface  $\Gamma$  we solve a PDE for  $\phi \in \mathbb{R}^d$  in the surface embedding domain  $\Omega_{\Gamma} \in \mathbb{R}^d$ :

$$\frac{\partial \phi}{\partial t} = F(\phi, \nabla \phi, \Delta \phi), \quad (2)$$

where  $\phi_{\Gamma}, \nabla_{\Gamma}, \Delta_{\Gamma}$  in (1) have been replaced by Cartesian operators in the embedding Euclidian space.

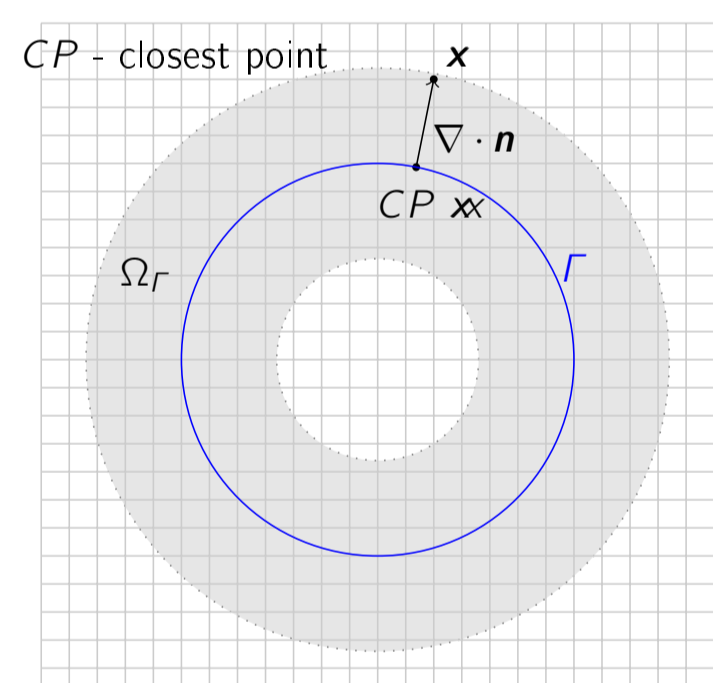


Fig 3: Closest point representation

**Solution procedure:**

1. Extend solution from surface  $\Gamma$  into the embedding space under the constraints  $\nabla u \cdot \mathbf{n} = 0$ ,  $u(\mathbf{x}_{\Gamma}) = u_{\Gamma}(\mathbf{x}_{\Gamma})$  via closest point method (Ruuth, Merriman, JCP 227, 2008), Fig. 3

2. Solve (2) in  $\Omega_{\Gamma}$  for timestep  $\Delta t$
3. Interpolate solution  $u(\mathbf{x})$  back onto the closest points (CP) on the surface
4. Repeat 1.-3. until end of simulation time

**Example:** Scalar advection on an ellipse:

$$\frac{\partial \phi_{\Gamma}}{\partial t} + \frac{\partial(\phi_{\Gamma} u_{\Gamma})}{\partial s} = 0$$

with arclength  $s$  and initial condition  $\phi(s, 0) = \cos^2(2\pi s/L)$ . The embedding PDE is

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{u}) = 0$$

where  $\phi$  and  $\mathbf{u}$  are closest point representations of  $\phi_{\Gamma}$  and  $u_{\Gamma}$ .

$\Delta x$	Error	Conv. rate
0.05	1.11e-03	-
0.025	2.76e-04	2.01
0.0125	6.90e-05	2.00
0.00625	1.80e-05	1.94
0.003125	4.29e-06	2.07

Tab. 1: Error in the  $L_{\infty}$ -norm using three-step Runge-Kutta time integration and 2nd order finite volume discretisation in space.

\* Kathrin Mellert, Diploma thesis, TU Berlin, to appear

## Sharp interface Poisson solver

We have implemented our locally second order Poisson solvers (JCP 2006, 2009) for node-centered and cell-centered problems into the LLNL AMR framework SAMRAI, Fig. 4. Key features

- finite volume discretization  $\int_{\partial\Omega} \beta (\nabla u \cdot \mathbf{n}) dS = \int_{\Omega^+} f^+ dV + \int_{\Omega^-} f^- dV - \int_{\Gamma} [\beta u_n] dS$ ,
- single piecewise bi- (2D) or tri-linear (3D) ansatz functions on normal cells,
- dual piecewise bi- (2D) or tri-linear (3D) ansatz functions on cut-cells with explicit incorporation of the known interface jump conditions  $[[u]]_{\Gamma}$  and  $[[\beta u_n]]_{\Gamma}$ ,
- 9- (2D) and 27-point (3D) compact stencils,
- locally second order accurate,
- interface representation using standard 2nd order levelsets.

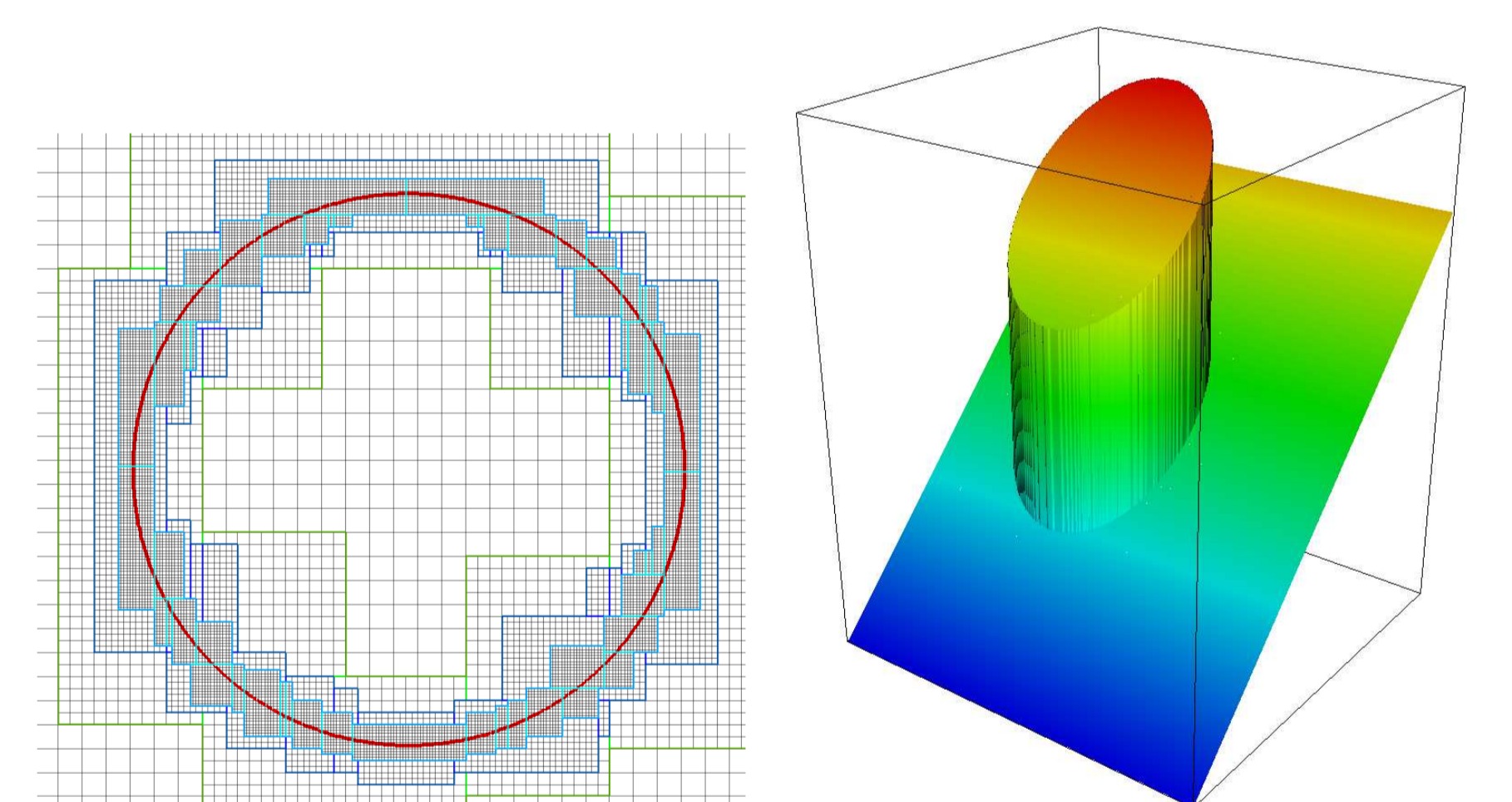


Fig. 4: Adapted Cartesian grid and initial hydrostatic pressure with density ratio 1.

## Asymptotics based sharp interface Poisson solver for arbitrary ratios of the coefficients

The accurate and efficient solution of elliptic equations

$$\nabla \cdot (\beta(\mathbf{x}) \nabla u(\mathbf{x})) = f(\mathbf{x})$$

with variable coefficients and prescribed interface jump conditions

$$[[u]]_{\Gamma} \quad \text{and} \quad [[\beta u_n]]_{\Gamma}$$

across an interface  $\Gamma$  is a key component in modeling incompressible two phase flow. One associated numerical problem is the robust solution in the case of a **large ratio of the coefficient** (e.g.  $\beta = \rho, \mu, D$ ).

**Asymptotic solution approach**

On each side of the interface we have

$$\begin{aligned} \nabla \cdot (\beta^+ \nabla u^+(\mathbf{x})) &= f^+(\mathbf{x}), \quad \mathbf{x} \in \Omega^+ \\ \nabla \cdot (\beta^- \nabla u^-(\mathbf{x})) &= f^-(\mathbf{x}), \quad \mathbf{x} \in \Omega^- \end{aligned} \quad (3)$$

with interface jump conditions

$$[[u]]_{\Gamma} = g(\mathbf{x}_{\Gamma}) \quad \text{and} \quad [[\beta u_n]]_{\Gamma} = h(\mathbf{x}_{\Gamma}).$$

We are interested in small ratios

$$\beta^- / \beta^+ = \varepsilon \ll 1$$

and write

$$\begin{aligned} \nabla \cdot (\varepsilon^{-1} \beta^- \nabla u^+(\mathbf{x})) &= f^+(\mathbf{x}), \quad \mathbf{x} \in \Omega^+ \\ \nabla \cdot (\beta^- \nabla u^-(\mathbf{x})) &= f^-(\mathbf{x}), \quad \mathbf{x} \in \Omega^- \end{aligned}$$

Introducing the asymptotic expansion

$$\begin{aligned} u^+(\mathbf{x}) &= u^{(0,+)}(\mathbf{x}) + \varepsilon u^{(1,+)}(\mathbf{x}) \\ u^-(\mathbf{x}) &= u^{(0,-)}(\mathbf{x}) + \varepsilon u^{(1,-)}(\mathbf{x}) \end{aligned}$$

into (3) leads to the following

**Solution strategy for the perturbation equations:**

1. Solve  $\nabla \cdot (\beta^- \nabla u^{(0,-)}(\mathbf{x})) = f^-(\mathbf{x})$

in  $\Omega^-$  with boundary condition  $u_n^{(0,-)} = u_{n,w}^-$  on  $\partial\Omega^- \setminus \Gamma$  and interface condition  $u^{(0,-)} = -g$  on  $\Gamma$ .

2. Solve the coupled problem

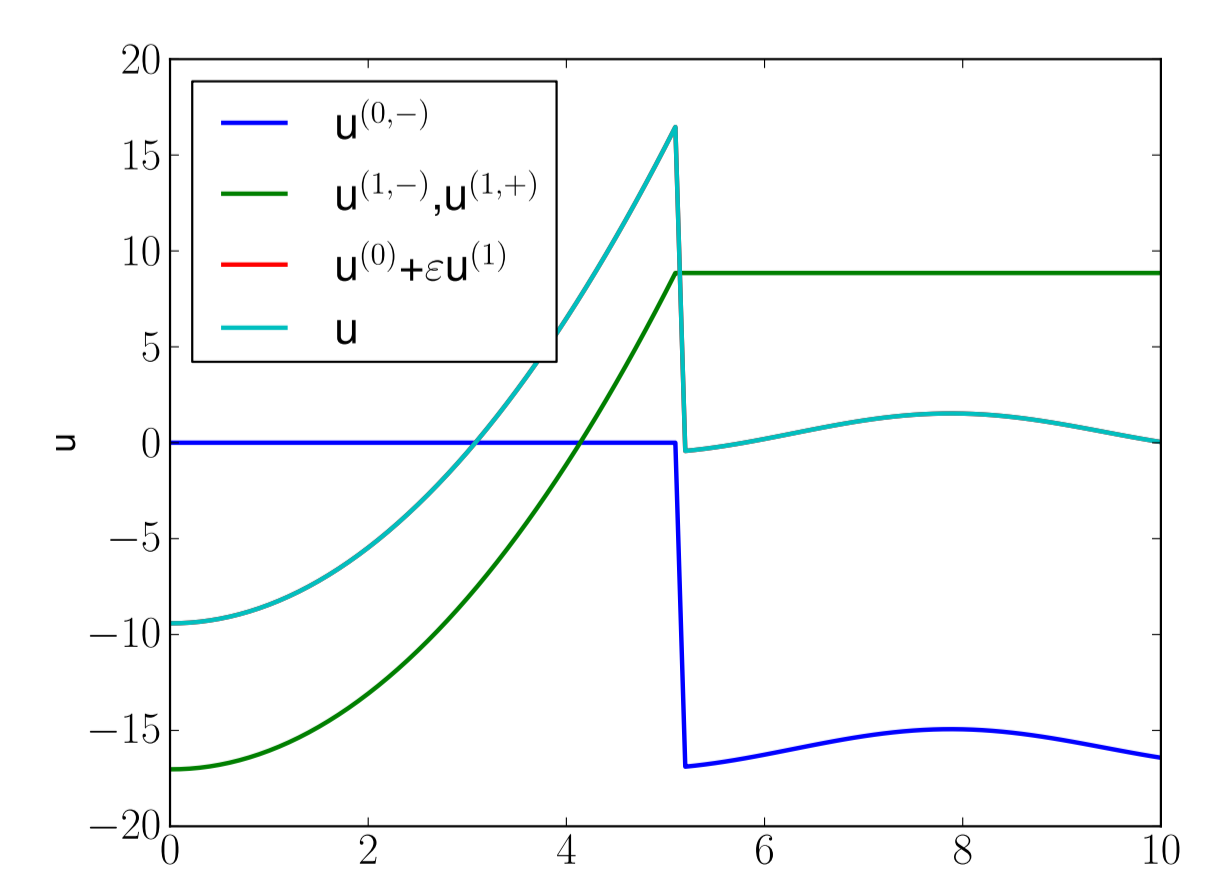
$$\begin{aligned} \nabla \cdot (\beta^- \nabla u^{(1,+)}(\mathbf{x})) &= f^+(\mathbf{x}), \quad \mathbf{x} \in \Omega^+ \\ \nabla \cdot (\beta^- \nabla u^{(1,-)}(\mathbf{x})) &= 0, \quad \mathbf{x} \in \Omega^- \end{aligned}$$

with

$$\begin{aligned} [[u^{(1)}]]_{\Gamma} &= u^{(1,+)} - u^{(1,-)} = 0, \\ \beta^- u_n^{(1,+)} - \varepsilon \beta^- u_n^{(1,-)} &= h + \beta^- u_n^{(0,-)} \quad \text{on } \Gamma, \\ u_n^{(1,-)} &= 0, \quad u_n^{(1,+)} = \varepsilon^{-1} u_{n,w}^+ \quad \text{on } \partial\Omega. \end{aligned}$$

**1D example:**

$$\begin{aligned} u^+ &= x^2, \quad x \in (0, 5.14] \\ u^- &= 10 + \sin(x), \quad x \in [5.14, 10) \end{aligned}$$



$\varepsilon = \beta^- / \beta^+$	$u^{(0,-)}$	$u^{(1)}$	$u$ (single step)	condition number
1	3.9e3	8.2e4	8.2e4	
1e-3	3.9e3	6.3e4	6.3e7	
1e-6	3.9e3	6.3e4	6.3e10	
1e-9	3.9e3	6.3e4	6.3e13	